

Regionalization Method for Nonlinear Differential Equation Systems In a Cartesian Plan

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Abstract: We propose a regionalization technique for analyzing nonlinear differential equation systems where coefficients are standard and nonzero. The present work starts with the study of a natural object which was the magic squares providing us with a new way to partition the plan in regions.

Key words: Internal Set Theory, Regionalization, nonlinear differential equations, operator

INTRODUCTION

In the present work we study the effects of delinearity of box type in orbital geometry and in orbital dynamics. We start this work with a first natural object around which we organize our plan , namely the magic square (MS) giving the partition of the plan \mathfrak{R}^2 into 25 external sets depicted in the following heuristic diagram .

15	22	9	16	3
2	14	21	8	20
19	1	13	25	7
6	18	5	12	24
23	10	17	4	11

Also, it plays an essential role in the nonclassic partition (regionalization) as following $(\mathfrak{R}, G_- \cup A_- \cup P \cup A_+ \cup G_+)$ of \mathfrak{R} .

Here are some of the questions that can be asked

- The apparition of the magic square
- The behaviour of the orbits in the regions of the MS
- The transition of the orbits of a region have the other
- The nature of the singular place
- The strut with the linear case.

The objects: We are interested by non-standard systems of nonlinear differential equations in \mathfrak{R}^2 provided with cartesian coordinates (X_1, X_2) .

$$\begin{cases} X_1' = a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} + b_1 \\ X_2' = a_{21} X_1^{[\alpha_1]} + a_{22} X_2^{[\alpha_2]} + b_2 \end{cases} \quad (1)$$

where the reals $a_{11}, a_{12}, a_{21}, a_{22}, b_1, b_2$ (resp $\alpha_1 > 0, \alpha_2 > 0$) are standard and nonzero (resp infinitely great)

$$\text{and for } i=1,2 \begin{cases} X_i^{[\alpha_i]} = X_i^{\alpha_i} & \text{if } X_i \geq 0 \\ X_i^{[\alpha_i]} = (-X_i)^{\alpha_i} & \text{if } X_i < 0 \end{cases}$$

we proceed to a sharp delinearization " of box type delinearization " of differential system with constant coefficients

$$\begin{cases} X_1' = a_{11} X_1 + a_{12} X_2 + b_1 \\ X_2' = a_{21} X_1 + a_{22} X_2 + b_2 \end{cases}$$

With conserving the linear equations. Then the problem is to evaluate the effects.

We can also see a problem of transient to the limit $(\alpha_1 \rightarrow +\infty, \alpha_2 \rightarrow +\infty)$

In some families which are linear when $\alpha_1 = \alpha_2 = 1$.

The differential system (1) is divided into three families : F_1, F_2, F_3

$$F_1: (a_{21}, a_{22}, b_2) = m(a_{11}, a_{12}, b_1)$$

$$F_2: (a_{21}, a_{22}, b_2) = m(a_{11}, a_{12}, b_1) + (0, 0, d); d \neq 0$$

$$F_3: \text{rang} \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} = 2$$

The technique: we use the technique of regionalization to get some predictions relative to the macroscopic behaviour of orbits and the dynamics along the orbits.

Some predictions

i. The magic square (MS) in the first analysis the plan \mathfrak{R}^2 is partitioned in 25 external regions where we denote by

(15) (resp (3)) the region defined by the conditions

$$X_1 \ll -1 \text{ and } X_2 \gg 1 \text{ (resp } X_1 \gg 1 \text{ and } X_2 \gg 1)$$

(22) (resp (16)) the region defined by the conditions

$X_1 \approx -1$ and $X_2 \succ \succ 1$ (resp $X_1 \approx 1$ and $X_2 \succ \succ 1$)
(9) (resp (17)) the region defined by the conditions

$X_1 \prec \prec -1$ and $X_2 \succ \succ 1$ (resp $|X_1| \prec \prec 1$ and $X_2 \prec \prec -1$)

(2) (resp (20)) the region defined by the conditions
 $X_1 \prec \prec -1$ and $X_2 \approx 1$ (resp $X_1 \succ \succ 1$ and $X_2 \approx 1$)

(14) (resp (8)) the region defined by the conditions
 $X_1 \approx -1$ and $X_2 \approx 1$ (resp $X_1 \approx 1$ and $X_2 \approx 1$)

(21) (resp (5)) the region defined by the conditions
 $|X_1| \prec \prec 1$ and $X_2 \approx 1$ (resp $|X_2| \prec \prec 1$ and $X_2 \approx -1$)

(19) (resp (7)) the region defined by the conditions
 $X_1 \prec \prec -1$ and $|X_2| \prec \prec 1$ (resp $X_1 \succ \succ 1$ and

$|X_2| \prec \prec 1$)

(1) (resp (25)) the region defined by the conditions
 $X_1 \approx -1$ and $|X_2| \prec \prec 1$ (resp $X_1 \approx 1$ and

$|X_2| \prec \prec 1$)

(1) (resp (25)) the region defined by the conditions
 $X_1 \prec \prec -1$ and $|X_2| \prec \prec 1$ (resp $X_1 \approx 1$ and

$|X_2| \prec \prec 1$)

(13) the region defined by the conditions

$|X_1| \prec \prec 1$ and $|X_2| \prec \prec 1$

(6) (resp (24)) the region defined by the conditions
 $X_1 \prec \prec -1$ and $X_2 \approx -1$ (resp $X_1 \succ \succ 1$ and

$X_2 \approx -1$)

(18) (resp (12)) the region defined by the conditions
 $X_1 \approx -1$ and $X_2 \approx -1$ (resp $X_1 \approx 1$ and $X_2 \approx -1$)

In (13) the macroscopic behaviour is determined by the rectiligne system

$$\begin{cases} X_1' \approx b_1 \\ X_2' \approx b_2 \end{cases}$$

In (9) (resp (17)) the macroscopic behaviour is determined by rectiligne system

$$\begin{cases} X_1' \approx a_{12} \\ X_2' \approx a_{22} \end{cases}$$

in (19)(resp (7)) the macroscopic behaviour is determined by the rectiligne system

$$\begin{cases} X_1' \approx a_{11} \\ X_2' \approx a_{21} \end{cases}$$

The refined magic square: to know the orbit behaviour in the other regions
(15,22,16,3,2,14,21,8,20,1,25,6,18,12,24,23,10,4,11)

We must refine the magic square (MS) by introducing the extensions 13',9',17',7',19' of the regions 13,9,17,7,19 respectively by the following conditions

$$13': a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} \in P \text{ and } a_{21} X_1^{[\alpha_1]} + a_{22} X_2^{[\alpha_2]} \in P$$

$$9' : X_2^{[\alpha_2]} \in G_+ \text{ and } \frac{X_1^{[\alpha_1]}}{X_2^{[\alpha_2]}} \in P$$

$$17' : X_2^{[\alpha_2]} \in G_- \text{ and } \frac{X_1^{[\alpha_1]}}{X_2^{[\alpha_2]}} \in P$$

$$7' : X_1^{[\alpha_1]} \in G_+ \text{ and } \frac{X_2^{[\alpha_2]}}{X_1^{[\alpha_1]}} \in P$$

$$19' : X_1^{[\alpha_1]} \in G_- \text{ and } \frac{X_2^{[\alpha_2]}}{X_1^{[\alpha_1]}} \in P$$

The family F₁: the function $F = X_2 - mX_1$ is a first integral.

The non-singular orbits are rectiligne and with the same slope m, the singular position is given by the equation

$$a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} + b_1 = 0$$

with the prime integral $F = X_2 - mX_1$ a differential system of the family F1 induce a family with one real parameter C (C specify the level straight) of differential equation of order 1 in \mathfrak{R} .

Namely, $D_C : a_{11} X_1^{[\alpha_1]} + a_{12} (mX_1 + C)^{[\alpha_2]} + b_1$

the form of singular place of F₁ implique some number of bifurcations in the equations D_c when C follow \mathfrak{R} .

The family F₂

$$F_2 \begin{cases} X_1' = a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} + b_1 \\ X_2' = m(a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} + b_1) + d, \quad d \neq 0 \end{cases}$$

The singular place is empty

If d was null we obtain a system of F₁ with a singular emptiness in the case to add $d \neq 0$ hunts the singular place, but it stay a witness that create a river phenomena.

The family F₃:
we distinguish three cases

j. the vectors $(a_{11}, a_{21}), (b_1, b_2), (a_{12}, a_{22})$ are pairwise independent.

$$jj. \begin{vmatrix} a_{11} & b_1 \\ a_{21} & b_2 \end{vmatrix} \neq 0 \quad ; \quad \begin{vmatrix} a_{12} & b_1 \\ a_{22} & b_2 \end{vmatrix} \neq 0$$

Proposition 1: In region (13) of magic square (MS) the macroscopic behaviour of orbits is determined by the rectiligne system

$$\begin{cases} X_1' \approx b_1 \\ X_2' \approx b_2 \end{cases}$$

Let $D = \frac{\partial}{\partial_i}$ denote the derivation operator as regard

X_i and state the following result :

Lemma 1: In the region (13) of the magic square (MS) the field of vectors

$Y = (a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1)D_1 + (a_{21}X_1^{[\alpha_1]} + a_{22}X_2^{[\alpha_2]} + b_2)D_2$
(associated to the system (1)) is infinitely near to the rectiligne field vectors $Y_{13} = b_1 D_1 + b_2 D_2$

Proof of lemma 1: consider the region (13) defined by the conditions

$$|X_1| \ll 1 \text{ and } |X_2| \ll 1$$

Or the field vectors:

$$Y = (a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1)D_1 + (a_{21}X_1^{[\alpha_1]} + a_{22}X_2^{[\alpha_2]} + b_2)D_2$$

Associated to the system

$$\begin{cases} X_1' = a_{11} X_1^{[\alpha_1]} + a_{12} X_2^{[\alpha_2]} + b_1 \\ X_2' = a_{21} X_1^{[\alpha_1]} + a_{22} X_2^{[\alpha_2]} + b_2 \end{cases} \quad (1)$$

as $|X_1| \ll 1$ and $|X_2| \ll 1$ and $\alpha_1 > 0, \alpha_2 > 0$ are real infinitude great

$X_1^{[\alpha_1]}$ and $X_2^{[\alpha_2]}$ are infinitude small.

consequently the field of vector Y in the coordinates (X_1, X_2) is infinitely near to the rectiligne field $Y_{13} = b_1 D_1 + b_2 D_2$

Proof of proposition 1: By the lemma of short shadows, the orbits of (1) have in the region (13) the same halo as the system orbits

$$\begin{cases} X_1' \approx b_1 \\ X_2' \approx b_2 \end{cases}$$

Proposition 2: In the regions (19) and (7) (resp (9) and (17)) of the magic square MS the macroscopic of the orbits behaviour has infinitely small fluctuations near is determined by the rectiligne system

$$\begin{cases} X_1' \approx a_{11} \\ X_2' \approx a_{21} \end{cases} \quad (\text{resp } \begin{cases} X_1' \approx a_{12} \\ X_2' \approx a_{22} \end{cases})$$

Lemma 2: In the regions (19) and (17) (resp 19' and (17')) of the magic square (MS) the field of vectors

$Y = (a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1)D_1 + (a_{21}X_1^{[\alpha_1]} + a_{22}X_2^{[\alpha_2]} + b_2)D_2$
Associated to the system (1) is infinitely near of the rectiligne field of vectors

$$Y_1 = \left(a_{11} + \frac{a_{12}X_2^{[\alpha_2]} + b_1}{X_1^{[\alpha_1]}} \right) D_1 + \left(a_{21} + \frac{a_{22}X_2^{[\alpha_2]} + b_2}{X_1^{[\alpha_1]}} \right) D_2$$

$$(\text{resp } Y_2 = \left(a_{12} + \frac{a_{11}X_1^{[\alpha_1]} + b_1}{X_2^{[\alpha_2]}} \right) D_1 + \left(a_{22} + \frac{a_{21}X_1^{[\alpha_1]} + b_2}{X_2^{[\alpha_2]}} \right) D_2)$$

and Y_1 (resp Y_2) is infinitely near of the rectiligne field of vectors

$$Y_1 = a_{11}D_1 + a_{21}D_2 \quad (\text{resp } Y_2 = a_{12}D_1 + a_{22}D_2).$$

Proof of lemma 2: In the regions (19) and (17)(resp (9) and (17)) of the magic square MS $X_2^{[\alpha_2]}$ (resp $X_1^{[\alpha_1]}$) is infinitely small $X_1^{[\alpha_1]}$ (resp $X_2^{[\alpha_2]}$) is infinitely great

since $\alpha_1 > 0, \alpha_2 > 0$ are infinitely great

Taking $X_1^{[\alpha_1]}$ (resp $X_2^{[\alpha_2]}$) as factor we can write the field of vectors :

$$Y = (a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1)D_1 + (a_{21}X_1^{[\alpha_1]} + a_{22}X_2^{[\alpha_2]} + b_2)D_2$$

Under the form

$$Y_1 = X_1^{[\alpha_1]} \left(a_{11} + \frac{a_{12}X_2^{[\alpha_2]} + b_1}{X_1^{[\alpha_1]}} \right) D_1 + \left(a_{21} + \frac{a_{22}X_2^{[\alpha_2]} + b_2}{X_2^{[\alpha_2]}} \right) D_2$$

$$(\text{resp } Y_2 = X_2^{[\alpha_2]} \left(a_{12} + \frac{a_{11}X_1^{[\alpha_1]} + b_1}{X_2^{[\alpha_2]}} \right) D_1 + \left(a_{22} + \frac{a_{21}X_1^{[\alpha_1]} + b_2}{X_2^{[\alpha_2]}} \right) D_2)$$

In the regions (19) and (7) (resp (9) and (17)) of the magic square the field Y has the same orbits as the field

$$Z_1 = \left(a_{11} + \frac{a_{12}X_2^{[\alpha_2]} + b_1}{X_1^{[\alpha_1]}} \right) D_1 + \left(a_{21} + \frac{a_{22}X_2^{[\alpha_2]} + b_2}{X_1^{[\alpha_1]}} \right) D_2$$

$$(\text{resp } Z_2 = \left(a_{12} + \frac{a_{11}X_1^{[\alpha_1]} + b_1}{X_2^{[\alpha_2]}} \right) D_1 + \left(a_{22} + \frac{a_{21}X_1^{[\alpha_1]} + b_2}{X_2^{[\alpha_2]}} \right) D_2)$$

The field Z_1 (resp Z_2) is infinitely near of the rectiligne field

$$Z_1' = a_{11}D_1 + a_{21}D_2 \quad (\text{resp } Z_2' = a_{12}D_1 + a_{22}D_2).$$

as the quantities

$$\left(\frac{a_{12}X_2^{[\alpha_2]} + b_1}{X_1^{[\alpha_1]}} \right) \quad (\text{resp } \left(\frac{a_{22}X_2^{[\alpha_2]} + b_2}{X_2^{[\alpha_2]}} \right))$$

$$\text{and } \left(\frac{a_{11}X_1^{[\alpha_1]} + b_1}{X_1^{[\alpha_1]}} \right) \quad (\text{resp } \left(\frac{a_{21}X_1^{[\alpha_1]} + b_2}{X_2^{[\alpha_2]}} \right))$$

are infinitely small .

Proof of proposition 2: The proposition 2 is an immediate consequence of the lemma of the short shadows as long as we have lemma 2.

Examination of the family F3:

$$F_3 : \text{rang} \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} = 2$$

We note

1. we can write

$$\begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} a_{11} \\ a_{21} \end{pmatrix} \alpha + \begin{pmatrix} a_{12} \\ a_{22} \end{pmatrix} \beta ; \quad \alpha^2 + \beta^2 > 0$$

$$\begin{cases} X_1' = a_{11}(X_1^{[\alpha_1]} + \alpha) + a_{12}(X_2^{[\alpha_2]} + \beta) \\ X_2' = a_{21}(X_1^{[\alpha_1]} + \alpha) + a_{22}(X_2^{[\alpha_2]} + \beta) \end{cases}$$

$$\alpha^2 + \beta^2 > 0$$

2. the singular point is reduced to the point S =

$$\left((-\alpha)^{\left[\frac{1}{\alpha_1}\right]}, (-\beta)^{\left[\frac{1}{\alpha_2}\right]} \right)$$

3. the singular point can only be find in the halo of a corner the unit square.

4. the slow region is given by $X_1^{[\alpha_1]} + \alpha \approx 0$, $X_2^{[\alpha_2]} + \beta \approx 0$ other

$$X_1 \in (\text{hal}(-\alpha))^{\left[\frac{1}{\alpha_1}\right]} \text{ and } X_2 \in (\text{hal}(-\beta))^{\left[\frac{1}{\alpha_2}\right]}$$

5. if $\alpha\beta \neq 0$, the slow region is strictly contained in the halo of $(\text{sign}(-\alpha), \text{sign}(-\beta))$.

if $\alpha = 0 \neq \beta$, the slow region is given by

$$X_1 \in P^{\left[\frac{1}{\alpha_1}\right]} \text{ and } X_2 \approx (\text{sign}(-\beta))^{\left[\frac{1}{\alpha_2}\right]}$$

if $\alpha \neq 0 = \beta$, the slow region is given by

$$X_2 \in P^{\left[\frac{1}{\alpha_2}\right]} \text{ and } X_1 \in (\text{hal}(-\alpha))^{\left[\frac{1}{\alpha_1}\right]}$$

6. the singular point is in the external square ES defined by $X_1^{[\alpha_1]} \in L$, $X_2^{[\alpha_2]} \in L$

Then in the region $X_1^{[\alpha_1]} \in G$, $X_2^{[\alpha_2]} \in G$, the curves $C_1 = (X_1' = 0)$ and $C_2 = (X_2' = 0)$ don't intersect.

Proposition 3

Given (p, q) a singular point of the system (1)

i. if $|p| \ll 1$ (resp $|q| \ll 1$) then:

* $|q| \approx 1$ (resp $|p| \approx 1$)

* $\frac{b_1}{b_2} \approx \frac{a_{12}}{a_{22}}$ (resp $\frac{b_1}{b_2} \approx \frac{a_{11}}{a_{21}}$)

ii. if $0 < |p| \ll 1$ (resp $0 < |q| \ll 1$) then

* $\frac{b_1}{b_2} \approx \frac{a_{12}}{a_{22}} \approx \frac{a_{11}}{a_{21}} \approx \beta$

* the function $F = X_1 - \beta X_2$ is a first integral of system (1)

* The singular place of the system (1) is defined by the equation $a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1 = 0$

iii. if $|p| > 1$ (resp $|q| > 1$) then

$\frac{q^{[\alpha_2]}}{p^{[\alpha_1]}}$ (resp $\frac{p^{[\alpha_1]}}{q^{[\alpha_2]}}$) is appreciated and we have $|q| > 1$ (resp $|p| > 1$)

Corollaire 1: a singular point of the system (1) taken in the region defined by $|X_1| > 1$ (resp $|X_2| > 1$)

And in the region A_{12} defined by the conditions

$$X_1 X_2 \neq 0 \text{ and } \frac{X_1^{[\alpha_1]}}{X_2^{[\alpha_2]}} \in A_+ \cup A_-$$

Proof of proposition 3

i. As (p, q) is not singular we have:

$$\begin{cases} a_{11} p^{[\alpha_1]} + a_{12} q^{[\alpha_2]} + b_1 = 0 \\ a_{21} p^{[\alpha_1]} + a_{22} q^{[\alpha_2]} + b_2 = 0 \end{cases}$$

if $|p| \ll 1$ (resp $|q| \ll 1$) then the real

$a_{12} p^{[\alpha_1]}$ and $a_{21} p^{[\alpha_1]}$ (resp $a_{12} q^{[\alpha_2]}$ and $a_{22} q^{[\alpha_2]}$) are infinitely small.

where the following relations

$$a_{12} q^{[\alpha_2]} + b_1 \approx 0 \text{ and } a_{22} q^{[\alpha_2]} + b_2 \approx 0$$

$$\text{(resp } a_{12} p^{[\alpha_1]} + b_1 \approx 0 \text{ and } a_{21} p^{[\alpha_1]} + b_2 \approx 0 \text{)}.$$

As a_{12} and a_{22} (resp a_{11} and a_{21}) are

appreciated we deduce the relations

$$q^{[\alpha_2]} \approx -\frac{b_1}{a_{12}} \approx -\frac{b_2}{a_{22}} \text{ (resp } p^{[\alpha_1]} \approx -\frac{b_1}{a_{11}} \approx -\frac{b_2}{a_{21}})$$

$$\text{but } \frac{b_1}{a_{12}} \text{ and } \frac{b_2}{a_{22}} \text{ (resp } \frac{b_1}{a_{11}} \text{ and } \frac{b_2}{a_{21}})$$

are standard.

$$\text{thus } \frac{b_1}{a_{12}} = \frac{b_2}{a_{22}} \text{ (resp } \frac{b_1}{a_{11}} = \frac{b_2}{a_{21}}).$$

$$\text{the relation } q^{[\alpha_2]} \approx -\frac{b_1}{a_{12}} \text{ (resp } p^{[\alpha_1]} \approx -\frac{b_1}{a_{11}})$$

show that

$$|q| \approx 1 \text{ (resp } |p| \approx 1) \text{ since } -\frac{b_1}{a_{12}} \text{ (resp } -\frac{b_1}{a_{11}}) \text{ is appreciated}$$

The point ii)

the obtained relation in (i)
 $\frac{b_1}{b_2} = \frac{a_{12}}{a_{22}} = \beta$ (resp $\frac{b_1}{b_2} = \frac{a_{11}}{a_{21}} = \beta$) implique the equality

$$a_{11}P^{[\alpha_1]} = -a_{12}q^{[\alpha_2]} - b_1 \quad (\text{resp } a_{22}q^{[\alpha_2]} = -a_{21}P^{[\alpha_1]} - b_2)$$

As

$$p \neq 0 \text{ (resp } q \neq 0) \text{ we obtain } a_{11} = \beta a_{22} \text{ (resp } \beta a_{22} = a_{12})$$

$$\text{or as } \beta = \frac{a_{11}}{a_{21}} = \frac{a_{22}}{a_{21}} \text{ (resp } \frac{a_{11}}{a_{22}} = \frac{a_{22}}{a_{21}})$$

$$\text{consequently } \frac{b_1}{b_2} = \frac{a_{22}}{a_{21}} = \frac{a_{22}}{a_{21}} = \beta$$

hence the equality

$$(a_{11}, a_{12}, b_1) = (a_{21}, a_{22}, b_2)$$

$$\text{Which imply the relation } X_1' - \beta X_2' = 0$$

So the function $F = X_1 - \beta X_2$ is a first integral of the system (1)

as $(a_{11}, a_{12}, b_1) = \beta (a_{21}, a_{22}, b_2)$ we see that

$$a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1 = \beta (a_{21}X_1^{[\alpha_1]} + a_{22}X_2^{[\alpha_2]} + b_2)$$

consequently the singular place of the system (1) is defined by the equation

$$a_{11}X_1^{[\alpha_1]} + a_{12}X_2^{[\alpha_2]} + b_1 = 0$$

The point iii)

If $|p| \gg 1$ (resp $|q| \gg 1$) the equality

$$a_{11}p^{[\alpha_1]} + a_{12}q^{[\alpha_2]} + b_1 = 0 \text{ can also be written}$$

$$a_{11} + a_{12} \frac{q^{[\alpha_2]}}{p^{[\alpha_1]}} + \frac{b_1}{p^{[\alpha_1]}} = 0 \text{ (resp } a_{12} + a_{11} \frac{p^{[\alpha_1]}}{q^{[\alpha_2]}} + \frac{b_1}{q^{[\alpha_2]}} = 0)$$

Show that $\frac{q^{[\alpha_2]}}{p^{[\alpha_1]}}$ (resp $\frac{p^{[\alpha_1]}}{q^{[\alpha_2]}}$) can not be infinitely

small or infinitely great

because a_{11} (resp a_{12}) is standard nonnull and

$$\frac{b_1}{p^{[\alpha_1]}} \text{ (resp } \frac{b_1}{q^{[\alpha_2]}}) \text{ infinitely small .}$$

Thus $\frac{q^{[\alpha_2]}}{p^{[\alpha_1]}}$ (resp $\frac{p^{[\alpha_1]}}{q^{[\alpha_2]}}$) is appreciated .

as $p^{[\alpha_1]}$ (resp $q^{[\alpha_2]}$) infinitely great, then it must be the same as $q^{[\alpha_2]}$ (resp $p^{[\alpha_1]}$) such that the quotient is appreciated .

if $q^{[\alpha_2]}$ (resp $p^{[\alpha_1]}$) is infinitely great , then $|q| \gg 1$ (resp $|p| \gg 1$) .

REFERENCES

1. Arnold, 1974. Equations Différentielle. Edition MIR. Moscou.
2. Bobo, S., 1988. Régionalisations et Van-der pol épaissi. Cahiers mathématiques d'Oran , fascicule n°2.
3. Bobo, S., 1988. Méthode des régionalisations. Quelques applications I.R.M.A Strasbourg.
4. Callot, J.L., 1981. Bifurcation du portrait de phase pour des équations différentielles du second ordre, Thèse Strasbourg.